

# Lattice dynamics and thermodynamics of light actinides

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## Abstract

In this paper, we present an *ab initio* study of the lattice dynamics of two light actinides, namely thorium and uranium. We worked within the framework of the quasi-harmonic approximation and by using density-functional theory, pseudo-potential method and linear response theory. We compare the phonon spectrums we obtained with inelastic neutron-scattering experiments.

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## 1. Introduction

Details about the phonon-dispersion relationships of actinides are very limited. For 30 years, experimental data were only available for Th and U. But recently, several works, using inelastic X-ray scattering, mostly on U [1,2] and Pu [3] have been published. Theoretical works are needed to tackle the f-electrons systems elastic properties. Unfortunately, these calculations are far from straightforward. The most important problem comes from the difficulty to treat correctly the f-electrons. To our knowledge, there is only one spectrum for a f-element beyond the semi-empirical methods, which was obtained by Dai et al. [4] using dynamical mean field theory (DMFT) on fcc plutonium.

In the harmonic approximation, the crystal free energy is built by adding to the 0 K isotherm a dynamical contribution which is approximated by the free energy of a system of harmonic oscillators. A step beyond the harmonic approximation can be made by allowing phonon frequencies to depend on crystal volume, the so-called quasi-harmonic approximation. Here, we apply the quasi-harmonic approximation to the study of the thermal properties of Th metal, such as thermal expansion, bulk modulus or Debye temperature. We show also the phonon spectrum of U.

## 2. Computational details

Our calculations are performed within the generalized gradient approximation (GGA) to density-functional theory, as im-

plemented in the ABINIT package [5,6]. Technical details on the computation of responses to atomic displacements can be found in Ref. [7] while Ref. [8] presents the subsequent computation of dynamical matrices and interatomic force constants.

The pseudo-potentials used in this work and their generation have been presented in Ref. [9]. They are norm conserving Trouillier and Martins [10] using GGA according to the recipe of Perdew et al. [11]. It has been shown that GGA represents a significant improvement on local density approximation (LDA) for f-elements. The kinetic energy cut-off of the plane wave bases is 160, 180 and 220 Ry for Th, U and Pu respectively. With these cut-offs, energy differences are converge within 1 mRy/at, and phonon frequencies within 0.01 THz. A Gaussian smearing of 0.02 Ry has been applied for Brillouin Zone integrations, which results in a metallic correction less than 0.1 mRy/at. In the electronic structure calculations, the exchange-correlation functional was consistent with the one used for the pseudo-potential generation.

## 3. Results and discussion

### 3.1. Thorium

We start with thorium which is the first 5f element with one electron f in the valence shell. The vibrational frequencies of Th were determined at several volumes within linear response framework. Dynamical matrices are computed at 19 wave ( $\mathbf{q}$ ) vectors in the irreducible wedge of BZ, i.e., on a  $8 \times 8 \times 8$   $\mathbf{q}$  grid, and are used for interpolation to obtain the bulk phonon dispersions. We present in Fig. 1, the phonons spectrum we ob-

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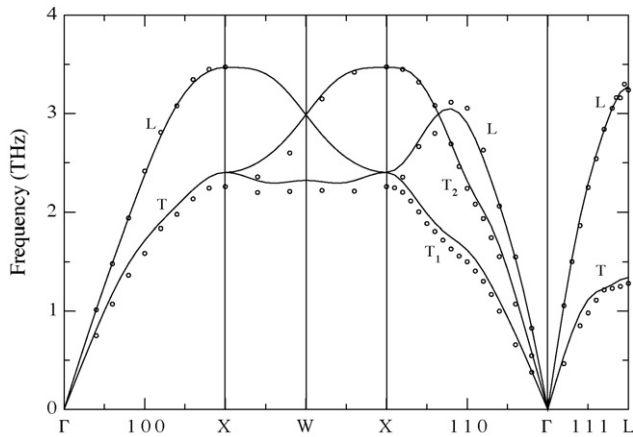


Fig. 1. Calculated phonon dispersion curves for fcc Th at the lattice parameter corresponding to static equilibrium. Experimental neutron-scattering data [12] are denoted by circles.

tained along several symmetry directions at zero pressure conditions. The neutron-scattering data obtained by Reese et al. [12] at room temperature are shown for comparison. The calculations reproduce really well the kinks observed in the  $T_1$  and  $T_2$  branches in the  $[110]$  direction. We have also calculated the phonons spectrum of Th in the bct structure. The results are presented in Fig. 2, and compared to the phonon dispersions of fcc. The  $T_1$  branch strongly softens during the phase transition in fcc and becomes imaginary (plotted with negative values) near the  $\Gamma$  point, in contrary, the frequencies of this branch increase with pressure in bct. This observation indicates that the fcc structure at this pressure is dynamically unstable.

Using the quasi-harmonic approximation, we have calculated the linear thermal expansion coefficient,  $\alpha_L$ . The value we obtained at room temperature,  $10 \times 10^{-6} \text{ K}^{-1}$  [13] is very closed to the experimental value of  $11 \times 10^{-6} \text{ K}^{-1}$  [14].

### 3.2. Uranium

The discovery of several charge density waves transition at low temperature in U is closely related to neutron-scattering experiments [15,16]. They showed that  $\alpha$ -U has very unusual dispersion curves in the  $[100]$  direction. In particular, the longitudinal optic-like mode,  $\Sigma_4$  shows a pronounced dip around  $[0.5, 0, 0]$ . The frequency of this mode decreases considerably with temperature [15,16]. At 70 K, it even contains an elastic component. Smith et al. concluded that the unit cell is doubled along the  $[100]$  direction and that a periodic lattice distortion has taken place.

In Fig. 3, we compare our results at 0 K to the neutron-scattering data obtained at room temperature. Let us mention that we found no influence of the spin-orbit coupling term on the phonon frequencies. We have followed the notation of Crummett et al. [15] to label the branches. As for thorium, the results are close to the experiments. The dip in the  $\Sigma_4$  mode is well reproduced. We found a more pronounced dip in the  $\Sigma_1$  branch than the one observed experimentally. This suggests that this branch is also strongly temperature-dependent. The motions associated with the lower  $\Sigma_1$  mode are primarily longitudinal acoustic-like

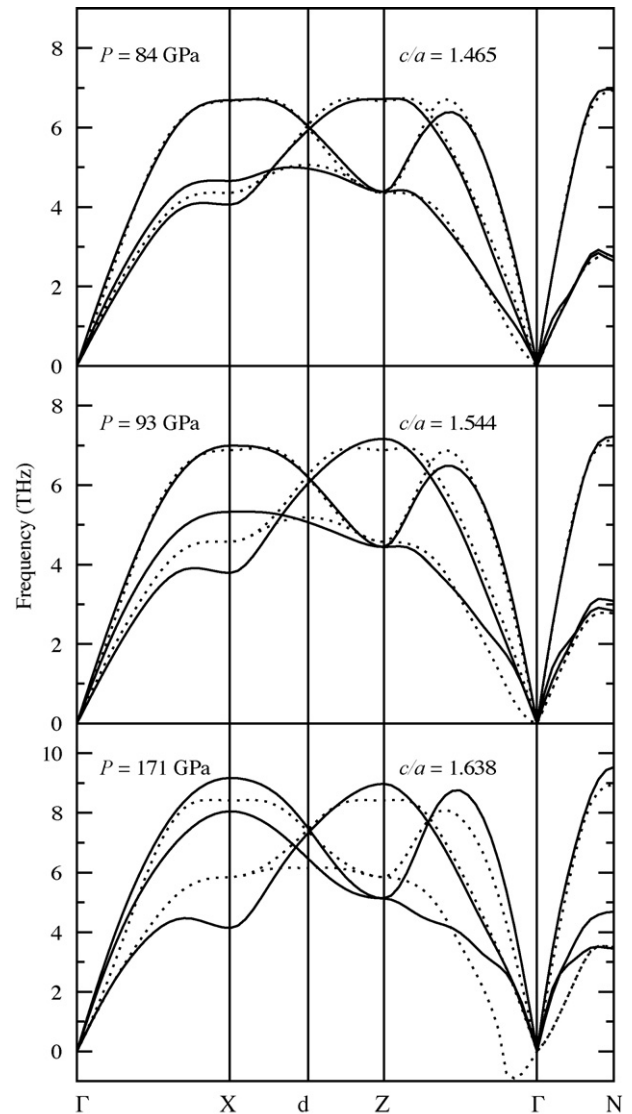


Fig. 2. Calculated phonon dispersion curves for bct (solid lines) and fcc (dashed lines) Th for pressures and  $c/a$  ratios in the transition region. The symmetry directions are those of fcc, the corresponding ones for bct are  $\Gamma$ -X-W-X- $\Gamma$ -L.

[15]. Therefore,  $\Sigma_4$  and  $\Sigma_1$  corresponds to out-of-phase and in-phase atomic displacements in the  $x$ -direction, respectively. The phase associated with the atomic motions involved in the CDW in the  $x$ -direction is  $(99 \pm 3)^\circ$  [17]. The motion is essentially optic ( $90^\circ$ ) but the acoustic part cannot be neglected and is reproduced in the softening observed in the  $\Sigma_1$  branch. The discrepancies we found between the theoretical and experimental spectrums are for branches related to an atomic motion in the  $x$ -direction, as  $\Delta_4$  and  $\Lambda_2$ . This is the direction where the largest atomic displacement occurs in the CDW, almost six times larger than the motions in  $y$ - and  $z$ -directions [18].

### 3.3. Conclusion

Phonon spectrum of the light actinides are now accessible with density-functional perturbation theory. We can now explore the phases diagrams of this class of materials in the framework of the quasi-harmonic approximation and obtain thermodynamic

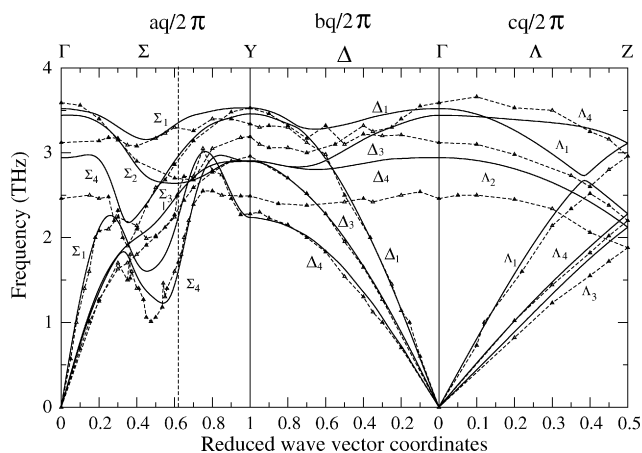


Fig. 3. Calculated phonon dispersion curves for  $\alpha$ -U at the lattice parameter corresponding to static equilibrium. The triangles are experimental neutron-scattering data [15].

quantities. This is very important since experiments are really difficult to realize in this class of materials.

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